

# The Cosmological Singularity

Marc Henneaux

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The first exactly solvable cosmological models of Einstein's theory revealed the presence of a very striking phenomenon : the Big Bang singularity (Friedmann, Lemaître).

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This raised the question as to whether this phenomenon was due to special simplifying assumptions or whether a singularity was a general property of the Einstein equations.

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This raised the question as to whether this phenomenon was due to special simplifying assumptions or whether a singularity was a general property of the Einstein equations.

Furthermore, if a singularity arises under general conditions, what is the analytical structure of this generic solution ?

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(other two : phase transitions and superconductivity)



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The BKL work showed that a singularity is a general property of a generic cosmological solution of the classical gravitational equations and not a consequence of the special symmetric structure of the exact models.

Most importantly, BKL were able to find the analytical structure of this generic solution and showed that its behaviour is of an extremely complex oscillatory character, of chaotic type.

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The asymptotic dynamics (as one goes toward the singularity) discovered by BKL can be understood in terms of a “cosmological billiard” system, where the cosmological evolution is described at each spatial point as the relativistic motion of a fictitious billiard ball moving in a region of hyperbolic space (Chitre, Misner).

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The billiard point of view provides a remarkable description of the gravitational field in the vicinity of a spacelike singularity.

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**The billiard point of view provides a remarkable description of the gravitational field in the vicinity of a spacelike singularity.**

In spite of the complexity of the Einstein-matter field equations, the asymptotic behaviour of the fields near a cosmological singularity can be phrased in surprisingly elementary terms involving finite-dimensional dynamical systems.

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**The billiard point of view provides a remarkable description of the gravitational field in the vicinity of a spacelike singularity.**

In spite of the complexity of the Einstein-matter field equations, the asymptotic behaviour of the fields near a cosmological singularity can be phrased in surprisingly elementary terms involving finite-dimensional dynamical systems.

**This description is valid generically, i.e. without making any symmetry assumption.**



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The billiard point of view has unexpectedly led to the discovery of a remarkable connection with one of the most beautiful and active subjects of modern mathematics, namely hyperbolic Coxeter groups and the theory of indefinite Kac Moody algebras (T. Damour-MH).

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This connection emerges because the billiard region in which the cosmological billiard ball moves turns out to possess exceptional properties,

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This intriguing fact opens up the fascinating perspective that an underlying infinite-dimensional symmetry algebra might play a central role in the fundamental formulation of gravity.

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These remarkable properties remain valid for string-related models.

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The purpose of the talk is to :



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- explain the main lines of the BKL analysis ;

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The purpose of the talk is to :

- explain the main lines of the BKL analysis ;
- discuss the properties of the “cosmological billiards” that emerge ;

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The purpose of the talk is to :

- explain the main lines of the BKL analysis ;
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- discuss the extension of the analysis to higher dimensions and matter-coupled models ;

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- explain the connection with fascinating mathematical structures ;

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- discuss the properties of the “cosmological billiards” that emerge ;
- discuss the extension of the analysis to higher dimensions and matter-coupled models ;
- explain the connection with fascinating mathematical structures ;
- briefly mention some open questions.

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Is a spacelike (“cosmological”) singularity generic? The way followed by BKL to answer this question was **constructive**.

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Assume that there is a spacelike singularity, and take it at  $t = 0$  for convenience.

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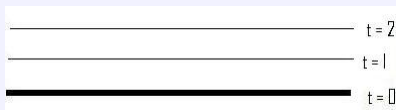
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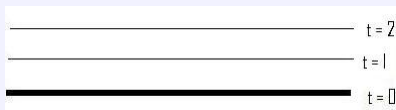
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Investigate then the Einstein equations as one approaches the singularity,  $t \rightarrow 0$ .

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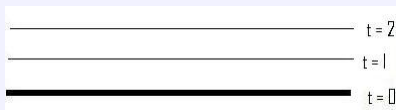
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BKL were able to find the general solution in the limit,

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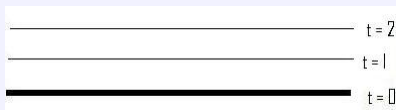
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BKL were able to find the general solution in the limit, and to show that this solution contained sufficiently many arbitrary physical functions to match generic initial data.

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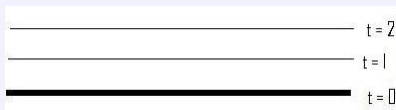
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Investigate then the Einstein equations as one approaches the singularity,  $t \rightarrow 0$ .

BKL were able to find the general solution in the limit, and to show that this solution contained sufficiently many arbitrary physical functions to match generic initial data.

This is remarkable given that the Einstein equations are non linear and remain so in the BKL limit.

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Their general solution can be described as follows.

Diagonalize metric (by “Iwasawa” change of frame -T. Damour, MH, H. Nicolai)

$$ds^2 = -dt^2 + a^2(t, \mathbf{x})\mathbf{l}^2 + b^2(t, \mathbf{x})\mathbf{m}^2 + c^2(t, \mathbf{x})\mathbf{n}^2$$

where  $a, b, c$  are the “scale factors” and  $\mathbf{l} = l_i(t, \mathbf{x}) dx^i$  etc

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Results : as  $t \rightarrow 0$

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Results : as  $t \rightarrow 0$

- the non-diagonal components freeze so that the interesting dynamics is carried by scale factors ;



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where  $a, b, c$  are the “scale factors” and  $\mathbf{l} = l_i(t, \mathbf{x})dx^i$  etc

Results : as  $t \rightarrow 0$

- the non-diagonal components freeze so that the interesting dynamics is carried by scale factors ;
- the dynamics of the scale factors decouples at each spatial point (equations become in the limit ODE’s with respect to time, a finite number at each point) - “ultralocality” ; not an assumption ! (“dimensional reduction without dimensional reduction”)

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The dynamics at each point of the scale factors follows the following rules.

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The dynamics at each point of the scale factors follows the following rules.

The spatial gradients can be neglected most of the time.

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The spatial gradients can be neglected most of the time.

When they are negligible, the scale factors have the “Kasner behaviour”  $a_i \sim t^{p_i}$ .

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Spatial gradients inevitably grow, however, and induce a “collision” from one Kasner regime to a new one. The collision is localized in time. After the collision has taken place, the spatial gradient term responsible for the collision decays, but another spatial gradient will then grow, leading to another collision.

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Terms due to the off-diagonal components also have a similar effect.

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When they are negligible, the scale factors have the “Kasner behaviour”  $a_i \sim t^{p_i}$ .

Spatial gradients inevitably grow, however, and induce a “collision” from one Kasner regime to a new one. The collision is localized in time. After the collision has taken place, the spatial gradient term responsible for the collision decays, but another spatial gradient will then grow, leading to another collision.

Terms due to the off-diagonal components also have a similar effect.

This leads to a billiard picture, where the Kasner regime corresponds to free flight motion, and where the effect of the spatial gradients and of the off-diagonal terms is to introduce “walls” against which there are collisions, changing one Kasner regime to another in a never-ending way.

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By appropriately parametrizing the scale factors, one can map this asymptotic billiard description into a region of 2-dimensional hyperbolic space  $\mathbb{H}_2$ .

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(2-dimensional because the scale factors are not independent due to the constraints).

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The Kasner motion is a geodesic on  $\mathbb{H}_2$ .

The walls are hyperplanes in  $\mathbb{H}_2$ .

There are two types of walls, “curvature walls” and “symmetry walls”.

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*The billiard table for four-dimensional pure gravity (one billiard system at each point in space) in the Poincaré disk model. Straight lines are symmetry walls. Arc of circles are gravitationnal walls. Thick lines indicate the dominant walls. The billiard table, which lies on the positive side of each wall, is indicated in grey. It forms a triangle with angles  $0$ ,  $\frac{\pi}{3}$ ,  $\frac{\pi}{2}$ .*

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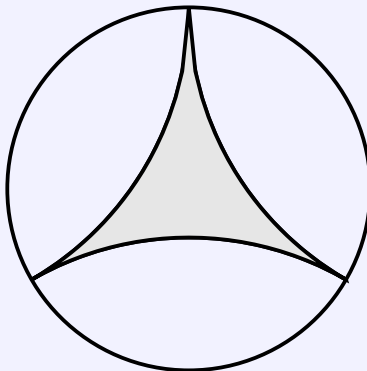
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*The billiard table for diagonal homogeneous cosmological models of Bianchi type IX. It is an ideal equilateral triangle formed by the three curvature (= gravitational) walls. The symmetry walls are absent. Any two different walls are parallel since they meet at infinity. The angles of this triangle are all equal to zero.*

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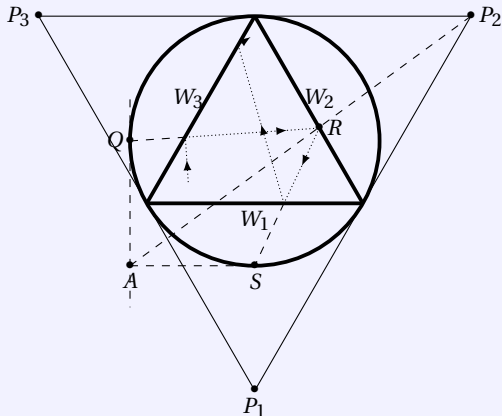
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*The motion is a succession of reflections in the walls  $W_1$ ,  $W_2$  and  $W_3$  and defines therefore an element in the group generated by the three reflections  $r_1$ ,  $r_2$ ,  $r_3$  in the three walls of the billiard table (drawn in the Klein model representation).*



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**As we have seen :**

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Chaos because of finite volume

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- $p$ -form fields play same role as non-diagonal degrees of freedom and curvature terms, and bring their own additional walls.

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- $p$ -form fields play same role as non-diagonal degrees of freedom and curvature terms, and bring their own additional walls.
- Dimension and shape of billiard table depends thus on spacetime dimension, menu of  $p$ -forms and dilaton couplings.
- Dynamics is chaotic if volume of billiard table is finite.



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Furthermore :

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Furthermore :

- The “billiard table” is a convex polyhedron in hyperbolic space bounded by hyperplanes  $H_s$  (“walls”), i.e., defined by

$$P = \bigcap_{s=1}^N H_s^+$$

where  $H_s^+$  is the positive half-space bounded by the hyperplane  $H_s$ , and  $N$  is the number of relevant bounding hyperplanes.

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where  $H_s^+$  is the positive half-space bounded by the hyperplane  $H_s$ , and  $N$  is the number of relevant bounding hyperplanes.

- There are four types of walls defining the billiard region : symmetry walls, gravitational walls,  $p$ -form electric walls,  $p$ -form magnetic walls.

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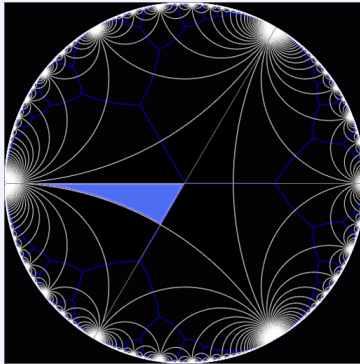
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## Angles between billiard walls



Angles between billiard walls :  $0 = \frac{\pi}{\infty}, \frac{\pi}{3}, \frac{\pi}{2}$ .

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For all string-related theories, the billiard table possesses the following remarkable properties :

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**It is acute-angled .**

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so that the product  $s_i s_j$  of the reflections  $s_i$  and  $s_j$ , which is a rotation by the angle  $\frac{2\pi}{m_{ij}}$ , fulfills

$$(s_i s_j)^{m_{ij}} = 1.$$

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A convex, acute-angled polyhedron with this property is a Coxeter polyhedron (here Coxeter simplex).

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**Because of these features,**

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A convex, acute-angled polyhedron with this property is a Coxeter polyhedron (here Coxeter simplex).

**Because of these features,**

- the billiard group is a Coxeter group.
- The billiard region is a fundamental domain for the action of the billiard group (= group generated by the reflections against the walls).

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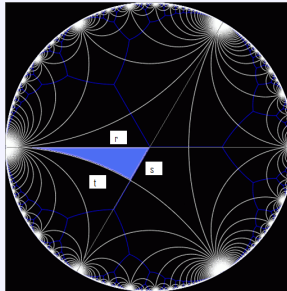
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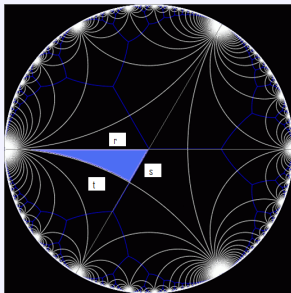
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## Coxeter group for pure gravity in $D = 4$



The reflections  $r$ ,  $s$  and  $t$ , subject to the Coxeter relations

$$s^2 = 1, r^2 = 1, t^2 = 1, (rs)^3 = (sr)^3 = 1, (ts)^2 = (st)^2 = 1,$$

generate the symmetry group (which is a Coxeter group). There are no other relations (standard Coxeter presentation) and the region in blue is a fundamental domain.

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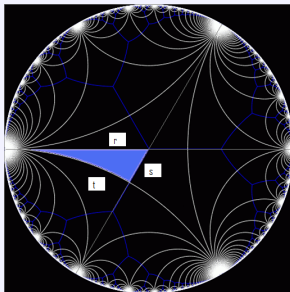
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## Coxeter group for pure gravity in $D = 4$



The reflections  $r$ ,  $s$  and  $t$ , subject to the Coxeter relations

$$s^2 = 1, r^2 = 1, t^2 = 1, (rs)^3 = (sr)^3 = 1, (ts)^2 = (st)^2 = 1,$$

generate the symmetry group (which is a Coxeter group). There are no other relations (standard Coxeter presentation) and the region in blue is a fundamental domain.

This is the group  $PGL(2, \mathbb{Z})$  (also noted  $A_1^{++}$ )

# Coxeter graphs

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One can associate a so-called “Coxeter graph” to Coxeter reflection groups. The Coxeter graph encodes the relations  $(s_i s_j)^{m_{ij}} = e$  among the generating reflections. If there is no relation among  $s_i$  and  $s_j$ , one conventionally declares  $m_{ij} = \infty$ . The graph contains as many vertices as generating reflections. For each pair of vertices, one draws an edge except when  $m_{ij} = 2$ , which corresponds to commuting reflections,  $s_i s_j = s_j s_i$ . Over each edge, one writes explicitly  $m_{ij}$ , except when  $m_{ij} = 3$  which is the “default value”.

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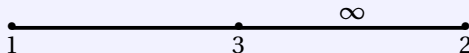
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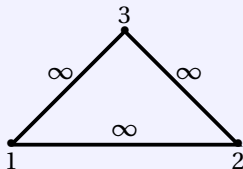
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In the case of the two reflection groups  $PGL(2, \mathbb{Z})$  and  $G\Gamma(2)$  (group of the ideal triangle), there are three vertices and the Coxeter graphs read respectively :



*The Coxeter graph for  $PGL(2, \mathbb{Z})$*



*The Coxeter graph for  $G\Gamma(2)$*

# Other theories

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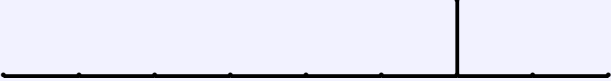
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Hyperbolic simplex Coxeter groups are very exceptional structures. The maximum possible rank is 10, and there are only 3 hyperbolic simplex Coxeter groups of rank 10 ...

Name	Coxeter graph
$B_8^{++} \equiv BE_{10}$	
$D_8^{++} \equiv DE_{10}$	
$E_8^{++} \equiv E_{10}$	

... and these are all realized by string-related supergravities!



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By definition, a Lorentzian Coxeter group is crystallographic if it stabilizes a lattice in the ambient Minkowski space  $\mathbb{R}^{M,1}$ .

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By definition, a Lorentzian Coxeter group is crystallographic if it stabilizes a lattice in the ambient Minkowski space  $\mathbb{R}^{M,1}$ .

A Coxeter group is crystallographic if and only if two conditions are satisfied : (i) The integers  $m_{ij}$  ( $i \neq j$ ) are restricted to be in the set  $\{2, 3, 4, 6, \infty\}$ , and (ii) for any closed circuit in the corresponding Coxeter graph, the number of edges labelled 4 or 6 is even.

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The Coxeter groups associated with string-related models are all simplex crystallographic Coxeter groups.

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Simplex, crystallographic Coxeter groups have a remarkable property : they can be identified with the Weyl groups of Kac-Moody algebras, traditionally denoted in the same way.

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Kac-Moody algebras are infinite-dimensional Lie algebras that generalize familiar finite-dimensional simple Lie algebras. The infinite-dimensional generalization preserves the key property of possessing a triangular decomposition such that the finite-dimensional concepts of Cartan subalgebra, raising and lowering operators, roots, positive roots, Weyl group, etc can all be introduced.

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Does this signal a hidden, infinite-dimensional symmetry that would exist independently of the cosmological context (the BKL analysis acting as a revelator) ?



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- Understanding spacelike singularities is a fundamental question.

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- Understanding spacelike singularities is a fundamental question.
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