

Effective quantum gravity from the point of view of perturbative algebraic QFT

Kasia Rejzner

University of York

Marcel Grossmann Meeting MG14, 13.07.2015



Outline of the talk



Quantum gravity

- Effective quantum gravity
- Symmetries
- Background independence



Algebraic quantum field theory

• A convenient framework to investigate conceptual problems in QFT is the Algebraic Quantum Field Theory (recently also perturbative AQFT).



Algebraic quantum field theory

- A convenient framework to investigate conceptual problems in QFT is the Algebraic Quantum Field Theory (recently also perturbative AQFT).
- It started as the axiomatic framework of Haag-Kastler [Haag & Kastler 64]: a model is defined by associating to each region 0 of Minkowski spacetime the algebra A(0) of observables that can be measured in 0.





Algebraic quantum field theory

- A convenient framework to investigate conceptual problems in QFT is the Algebraic Quantum Field Theory (recently also perturbative AQFT).
- It started as the axiomatic framework of Haag-Kastler [Haag & Kastler 64]: a model is defined by associating to each region (0) of Minkowski spacetime the algebra $\mathfrak{A}(0)$ of observables that can be measured in (0).
- The physical notion of subsystems is realized by the condition of isotony, i.e.: O₂ ⊃ O₁ ⇒ A(O₂) ⊃ A(O₁). We obtain a net of algebras.



Different aspects of AQFT and relations to physics





Difficulties in QFT on curved spacetimes

To include effects of general relativity into QFT, one has to be able to describe quantum fields on a general class of spacetimes.





To include effects of general relativity into QFT, one has to be able to describe quantum fields on a general class of spacetimes.

LCOFT

• Generically, the group of spacetime symmetries is trivial, so concept of particles as irreducible representations of such a group doesn't make sense.





To include effects of general relativity into QFT, one has to be able to describe quantum fields on a general class of spacetimes.

- Generically, the group of spacetime symmetries is trivial, so concept of particles as irreducible representations of such a group doesn't make sense.
- The concept of the vacuum as the state with no particles also becomes meaningless.





To include effects of general relativity into QFT, one has to be able to describe quantum fields on a general class of spacetimes.

- Generically, the group of spacetime symmetries is trivial, so concept of particles as irreducible representations of such a group doesn't make sense.
- The concept of the vacuum as the state with no particles also becomes meaningless.
- Transition to imaginary times (Wick rotation) is possible only in special cases.





To include effects of general relativity into QFT, one has to be able to describe quantum fields on a general class of spacetimes.

- Generically, the group of spacetime symmetries is trivial, so concept of particles as irreducible representations of such a group doesn't make sense.
- The concept of the vacuum as the state with no particles also becomes meaningless.
- Transition to imaginary times (Wick rotation) is possible only in special cases.
- Problems with the Fourier transform: calculations relying on momentum space cannot be performed.





• These conceptual problems can be easily solved in the algebraic approach.





- These conceptual problems can be easily solved in the algebraic approach.
- The corresponding generalization of AQFT is called locally covariant quantum field theory [Hollands & Wald CMP 01 Brunetti, Fredenhagen & Verch CMP 01, Fewster & Verch AHP 12,...],



• These conceptual problems can be easily solved in the algebraic approach.

- The corresponding generalization of AQFT is called locally covariant quantum field theory [Hollands & Wald CMP 01 Brunetti, Fredenhagen & Verch CMP 01, Fewster & Verch AHP 12,...],
- Come to the QFT parallel sessions!



• These conceptual problems can be easily solved in the algebraic approach.

LCOFT

- The corresponding generalization of AQFT is called locally covariant quantum field theory [Hollands & Wald CMP 01 Brunetti, Fredenhagen & Verch CMP 01, Fewster & Verch AHP 12,...],
- Come to the QFT parallel sessions!

Main advantages





• These conceptual problems can be easily solved in the algebraic approach.

LCOFT

- The corresponding generalization of AQFT is called locally covariant quantum field theory [Hollands & Wald CMP 01 Brunetti, Fredenhagen & Verch CMP 01, Fewster & Verch AHP 12,...],
- Come to the QFT parallel sessions!

Main advantages

• Local algebras of observables $\mathfrak{A}(\mathfrak{O})$ are defined abstractly, the Hilbert space representation comes later (this deals with the non-uniqueness of the vacuum).



• These conceptual problems can be easily solved in the algebraic approach.

LCOFT

- The corresponding generalization of AQFT is called locally covariant quantum field theory [Hollands & Wald CMP 01 Brunetti, Fredenhagen & Verch CMP 01, Fewster & Verch AHP 12,...],
- Come to the QFT parallel sessions!

Main advantages

- Local algebras of observables $\mathfrak{A}(\mathfrak{O})$ are defined abstractly, the Hilbert space representation comes later (this deals with the non-uniqueness of the vacuum).
- $\bullet\,$ Algebras $\mathfrak{A}(\mathfrak{O})$ are constructed using only the local data.



• These conceptual problems can be easily solved in the algebraic approach.

LCOFT

- The corresponding generalization of AQFT is called locally covariant quantum field theory [Hollands & Wald CMP 01 Brunetti, Fredenhagen & Verch CMP 01, Fewster & Verch AHP 12,...],
- Come to the QFT parallel sessions!

Main advantages

- Local algebras of observables $\mathfrak{A}(\mathfrak{O})$ are defined abstractly, the Hilbert space representation comes later (this deals with the non-uniqueness of the vacuum).
- $\bullet\,$ Algebras $\mathfrak{A}(\mathfrak{O})$ are constructed using only the local data.
- Local features of the theory (observables) are separated from the global features (states).



LCOFT

• In the original AQFT axioms we associate algebras to regions of a fixed spacetimes. Now we go a step further.

- In the original AQFT axioms we associate algebras to regions of a fixed spacetimes. Now we go a step further.
- Replace O₁ and O₂ with arbitrary spacetimes *M* = (*M*, *g*), *N* = (*N*, *g'*) and require the embedding ψ : *M* → *N* to be an isometry.



- In the original AQFT axioms we associate algebras to regions of a fixed spacetimes. Now we go a step further.
- Replace O₁ and O₂ with arbitrary spacetimes *M* = (*M*, *g*), *N* = (*N*, *g'*) and require the embedding ψ : *M* → *N* to be an isometry.
- Require that ψ preserves orientations and the causal structure (no new causal links are created by the embedding).



- In the original AQFT axioms we associate algebras to regions of a fixed spacetimes. Now we go a step further.
- Replace O₁ and O₂ with arbitrary spacetimes M = (M, g), N = (N, g') and require the embedding ψ : M → N to be an isometry.
- Assign to each spacetime M an algebra *A*(M) and to each admissible embedding ψ a homomorphism of algebras *Aψ* (notion of subsystems). This has to be done covariantly.



AQFT LCQFT

Locally covariant fields

• In the framework of LCQFT, locally covariant fields are used to identify (put labels on) observables localized in different region of spacetime, in the absence of symmetries.





Locally covariant fields

- In the framework of LCQFT, locally covariant fields are used to identify (put labels on) observables localized in different region of spacetime, in the absence of symmetries.
- Let D(O) denote the space of test functions supported in O. An LC field is a family of maps Φ_M : D(M) → A(M), labeled by spacetimes M such that:
 Aψ(Φ_O(f)) = Φ_M(ψ_{*}f)[h].





Locally covariant fields

- In the framework of LCQFT, locally covariant fields are used to identify (put labels on) observables localized in different region of spacetime, in the absence of symmetries.
- Let D(O) denote the space of test functions supported in O. An LC field is a family of maps Φ_M : D(M) → A(M), labeled by spacetimes M such that:
 Aψ(Φ_O(f)) = Φ_M(ψ_{*}f)[h].
- This generalizes the notion of Wightman's operator-valued distributions





Locally covariant fields

- In the framework of LCQFT, locally covariant fields are used to identify (put labels on) observables localized in different region of spacetime, in the absence of symmetries.
- Let D(O) denote the space of test functions supported in O. An LC field is a family of maps Φ_M : D(M) → A(M), labeled by spacetimes M such that:
 Aψ(Φ_O(f)) = Φ_M(ψ_{*}f)[h].
- This generalizes the notion of Wightman's operator-valued distributions
- Locally covariant fields are candidates for observables in GR.



Effective quantum gravity Symmetries Background independence



Difficulties in quantum gravity

• In contrast to QFT on curved spacetimes, in QG the spacetime structure is dynamical. Need for background independence.



Effective quantum gravity Symmetries Background independence



Difficulties in quantum gravity

- In contrast to QFT on curved spacetimes, in QG the spacetime structure is dynamical. Need for background independence.
- "Points" lose their meaning. The theory is invariant under diffeomorphism transformations.



Effective quantum gravity Symmetries Background independence



Difficulties in quantum gravity

- In contrast to QFT on curved spacetimes, in QG the spacetime structure is dynamical. Need for background independence.
- "Points" lose their meaning. The theory is invariant under diffeomorphism transformations.
- As a QFT, quantum gravity is power counting non-renormalizable.



Effective quantum gravity Symmetries Background independence



Ways around some of the problems

Based on our recent work: R. Brunetti, K. Fredenhagen, KR, *Quantum gravity from the point of view of locally covariant quantum field theory*, [arXiv:1306.1058].





Ways around some of the problems

Based on our recent work: R. Brunetti, K. Fredenhagen, KR, *Quantum gravity from the point of view of locally covariant quantum field theory*, [arXiv:1306.1058].



• Non-renormalizability: use Epstein-Glaser renormalization to obtain finite results for any fixed energy scale. Think of the theory as an effective theory.



Ways around some of the problems

Based on our recent work: R. Brunetti, K. Fredenhagen, KR, *Quantum gravity from the point of view of locally covariant quantum field theory*, [arXiv:1306.1058].



- Non-renormalizability: use Epstein-Glaser renormalization to obtain finite results for any fixed energy scale. Think of the theory as an effective theory.
- **Dynamical nature of spacetime**: make a split of the metric into background and perturbation, quantize the perturbation as a quantum field on a curved background, show background independence at the end.



Ways around some of the problems

Based on our recent work: R. Brunetti, K. Fredenhagen, KR, *Quantum gravity from the point of view of locally covariant quantum field theory*, [arXiv:1306.1058].



- Non-renormalizability: use Epstein-Glaser renormalization to obtain finite results for any fixed energy scale. Think of the theory as an effective theory.
- **Dynamical nature of spacetime**: make a split of the metric into background and perturbation, quantize the perturbation as a quantum field on a curved background, show background independence at the end.
- **Diffeomorphism invariance**: use the BV formalism to do the gauge fixing. Possible difficulties: base manifold is Lorentzian and non-compact, symmetry group is infinite dimensional, so is the space of metrics.

Effective quantum gravity Symmetries Background independence



• In experiment, geometric structure is probed by local observations. We have the following data:

Effective quantum gravity Symmetries Background independence

- In experiment, geometric structure is probed by local observations. We have the following data:
 - Some region O of spacetime where the measurement is performed,



Effective quantum gravity Symmetries Background independence

- In experiment, geometric structure is probed by local observations. We have the following data:
 - Some region O of spacetime where the measurement is performed,
 - An observable Φ , which we measure,



Effective quantum gravity Symmetries Background independence

- In experiment, geometric structure is probed by local observations. We have the following data:
 - Some region O of spacetime where the measurement is performed,
 - An observable Φ , which we measure,
 - We don't measure the observable (e.g. curvature) at a point. This is modeled by smearing with a test function *f*. For example:

$$\Phi(f) = \int f(x)R(x)d\mu(x).$$



Effective quantum gravity Symmetries Background independence

Intuitive idea

- In experiment, geometric structure is probed by local observations. We have the following data:
 - Some region O of spacetime where the measurement is performed,
 - An observable Φ , which we measure,
 - We don't measure the observable (e.g. curvature) at a point. This is modeled by smearing with a test function *f*. For example:

$$\Phi(f) = \int f(x)R(x)d\mu(x).$$

• Think of the measured observable as a function of a perturbation of the fixed background metric: $\tilde{g}_{\mu\nu} = g_{\mu\nu} + h_{\mu\nu}.$

Effective quantum gravity Symmetries Background independence

- In experiment, geometric structure is probed by local observations. We have the following data:
 - Some region O of spacetime where the measurement is performed,
 - An observable Φ , which we measure,
 - We don't measure the observable (e.g. curvature) at a point. This is modeled by smearing with a test function *f*. For example:

$$\Phi(f) = \int f(x)R(x)d\mu(x).$$

- Think of the measured observable as a function of a perturbation of the fixed background metric: $\tilde{g}_{\mu\nu} = g_{\mu\nu} + h_{\mu\nu}$.
- Diffeomorphism transformation: move our experimental setup to a different region O'.





Effective quantum gravity Symmetries Background independence



Building models in LCQFT

• One of the methods to build models in LCQFT is the so called functional approach.



- One of the methods to build models in LCQFT is the so called functional approach.
- The main idea is to model observables as functionals on the the space 𝔅(𝔐) of off-shell field configurations. For the effective theory of gravity 𝔅(𝔐) = Γ((T*M)^{⊗2}).



- One of the methods to build models in LCQFT is the so called functional approach.
- The main idea is to model observables as functionals on the the space 𝔅(𝔐) of off-shell field configurations. For the effective theory of gravity 𝔅(𝔐) = Γ((T*M)^{⊗2}).
- On the space $\mathfrak{F}(\mathcal{M})$ of sufficiently well behaving functionals we introduce first the classical dynamics by defining a Poisson structure.



- One of the methods to build models in LCQFT is the so called functional approach.
- The main idea is to model observables as functionals on the the space 𝔅(𝔐) of off-shell field configurations. For the effective theory of gravity 𝔅(𝔐) = Γ((T*M)^{⊗2}).
- On the space $\mathfrak{F}(\mathcal{M})$ of sufficiently well behaving functionals we introduce first the classical dynamics by defining a Poisson structure.
- Next, we use the deformation quantization to construct the non-commutative quantum algebra.



- One of the methods to build models in LCQFT is the so called functional approach.
- The main idea is to model observables as functionals on the the space 𝔅(𝔐) of off-shell field configurations. For the effective theory of gravity 𝔅(𝔐) = Γ((T*M)^{⊗2}).
- On the space $\mathfrak{F}(\mathcal{M})$ of sufficiently well behaving functionals we introduce first the classical dynamics by defining a Poisson structure.
- Next, we use the deformation quantization to construct the non-commutative quantum algebra.
- We work all the time on the same set of functionals, but we equip it with different algebraic structures (i.e. Poisson bracket, non-commutative * product).

Effective quantum gravity Symmetries Background independence



Diffeomorphism invariance

We consider an LC field Φ as a family of maps
 Φ_M : D(M) → 𝔅(M), where M ≡ (M, g).



Diffeomorphism invariance

- We consider an LC field Φ as a family of maps $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$, where $\mathcal{M} \equiv (M, g)$.
- Let $\xi \in \Gamma(TM)$ be an infinitesimal diffeomorphism. It acts on $\Phi_{(M,g)}(f)$ as

 $(\rho(\xi)\Phi)_{(M,g)}(f)[h] =$ $\left\langle (\Phi_{(M,g)}(f))^{(1)}[h], \pounds_{\xi}(g+h) \right\rangle + \Phi_{(M,g)}(\pounds_{\xi}f)[h]$



- We consider an LC field Φ as a family of maps $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$, where $\mathcal{M} \equiv (M, g)$.
- Let $\xi \in \Gamma(TM)$ be an infinitesimal diffeomorphism. It acts on $\Phi_{(M,g)}(f)$ as

 $(\rho(\xi)\Phi)_{(M,g)}(f)[h] =$ $\left\langle (\Phi_{(M,g)}(f))^{(1)}[h], \pounds_{\xi}(g+h) \right\rangle + \Phi_{(M,g)}(\pounds_{\xi}f)[h]$

• Infinitesimal diffeomorphism invariance is the condition that: $\rho(\xi)\Phi = 0.$



- We consider an LC field Φ as a family of maps $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$, where $\mathcal{M} \equiv (M, g)$.
- Let $\xi \in \Gamma(TM)$ be an infinitesimal diffeomorphism. It acts on $\Phi_{(M,g)}(f)$ as

 $(\rho(\xi)\Phi)_{(M,g)}(f)[h] = \left\langle (\Phi_{(M,g)}(f))^{(1)}[h], \pounds_{\xi}(g+h) \right\rangle + \Phi_{(M,g)}(\pounds_{\xi}f)[h]$

- Infinitesimal diffeomorphism invariance is the condition that: $\rho(\xi)\Phi = 0.$
- Example: $\int R[g+h] f d\mu_{g+h}$ is diffeomorphism invariant, but $\int R[g+h] f d\mu_g$ is not.

Effective quantum gravity Symmetries Background independence



Relational observables

 We realize the choice of a coordinate system by introducing four scalar fields X^μ, which parametrize points of spacetime.



Relational observables

- We realize the choice of a coordinate system by introducing four scalar fields X^μ, which parametrize points of spacetime.
- Fix a function *f* : ℝ⁴ → ℝ, then the change of *f* = *X***f* due to the change of the coordinate system is realized through the change of scalar fields *X^µ*.



- We realize the choice of a coordinate system by introducing four scalar fields X^μ, which parametrize points of spacetime.
- Fix a function *f* : ℝ⁴ → ℝ, then the change of *f* = *X***f* due to the change of the coordinate system is realized through the change of scalar fields *X^µ*.
- For an LC field Φ we obtain a map

$$\Phi_f(h,X) \doteq \Phi_{(M,g)}(X^*f)(h) \,,$$



- We realize the choice of a coordinate system by introducing four scalar fields X^μ, which parametrize points of spacetime.
- Fix a function *f* : ℝ⁴ → ℝ, then the change of *f* = *X***f* due to the change of the coordinate system is realized through the change of scalar fields *X^μ*.
- For an LC field Φ we obtain a map

$$\Phi_f(h,X) \doteq \Phi_{(M,g)}(X^*f)(h) \,,$$

• In the next step make X^{μ} dynamical either by interpreting them as 4 dust fields (Brown-Kuchař model) or by constructing them locally from the metric and its derivatives.



- We realize the choice of a coordinate system by introducing four scalar fields X^μ, which parametrize points of spacetime.
- Fix a function *f* : ℝ⁴ → ℝ, then the change of *f* = *X***f* due to the change of the coordinate system is realized through the change of scalar fields *X^μ*.
- For an LC field Φ we obtain a map

$$\Phi_f(h,X) \doteq \Phi_{(M,g)}(X^*f)(h) \,,$$

- In the next step make X^{μ} dynamical either by interpreting them as 4 dust fields (Brown-Kuchař model) or by constructing them locally from the metric and its derivatives.
- Denote such metric dependent coordinates by X_{g+h}^{μ} .



- We realize the choice of a coordinate system by introducing four scalar fields X^μ, which parametrize points of spacetime.
- Fix a function *f* : ℝ⁴ → ℝ, then the change of *f* = *X***f* due to the change of the coordinate system is realized through the change of scalar fields *X^µ*.
- For an LC field Φ we obtain a map

$$\Phi_f(h,X) \doteq \Phi_{(M,g)}(X^*f)(h) \,,$$

- In the next step make X^{μ} dynamical either by interpreting them as 4 dust fields (Brown-Kuchař model) or by constructing them locally from the metric and its derivatives.
- Denote such metric dependent coordinates by X_{g+h}^{μ} .
- Each Φ_f induces a relational observable $g \mapsto \Phi_f(h, X_{g+h})$.

Effective quantum gravity Symmetries Background independence

Relative Cauchy evolution

 Let N₊ and N₋ be two spacetimes that embed into two other spacetimes M₁ and M₂ around Cauchy surfaces, via admissible embeddings χ_{k,±}, k = 1, 2.



Relative Cauchy evolution

 Let N₊ and N₋ be two spacetimes that embed into two other spacetimes M₁ and M₂ around Cauchy surfaces, via admissible embeddings χ_{k,±}, k = 1, 2.

• Then

 $\beta = \mathfrak{A}\chi_{1+} \circ (\mathfrak{A}\chi_{2+})^{-1} \circ \mathfrak{A}\chi_{2-} \circ (\mathfrak{A}\chi_{1-})^{-1}$ is an automorphism of $\mathfrak{A}(\mathcal{M}_1)$. This is the consequence of the Time-slice axiom of LCQFT.



Effective quantum gravity Symmetries Background independence



Background independence

• Let $\mathcal{M}_1 = (M, g_1)$ and $\mathcal{M}_2 = (M, g_2)$, where g_1 and g_2 differ by a (compactly supported) symmetric tensor *h* (see the diagram).



Background independence

• Let $\mathcal{M}_1 = (M, g_1)$ and $\mathcal{M}_2 = (M, g_2)$, where g_1 and g_2 differ by a (compactly supported) symmetric tensor *h* (see the diagram).

• Define
$$\Theta_{\mu\nu}(x) \doteq \frac{\delta\beta_h}{\delta h_{\mu\nu}(x)}\Big|_{h=0}$$
.



Background independence

• Let $\mathcal{M}_1 = (M, g_1)$ and $\mathcal{M}_2 = (M, g_2)$, where g_1 and g_2 differ by a (compactly supported) symmetric tensor *h* (see the diagram).

• Define
$$\Theta_{\mu\nu}(x) \doteq \frac{\delta\beta_h}{\delta h_{\mu\nu}(x)}\Big|_{h=0}$$
.

 The infinitesimal background independence is the condition Θ_{μν} = 0.



Background independence

• Let $\mathcal{M}_1 = (M, g_1)$ and $\mathcal{M}_2 = (M, g_2)$, where g_1 and g_2 differ by a (compactly supported) symmetric tensor *h* (see the diagram).

• Define
$$\Theta_{\mu\nu}(x) \doteq \frac{\delta\beta_h}{\delta h_{\mu\nu}(x)}\Big|_{h=0}$$
.

- The infinitesimal background independence is the condition Θ_{µν} = 0.
- We have proven that this condition is fulfilled as a consequence of quantized Einstein's equations.





• AQFT is a convenient framework to solve conceptual problems of QFT on curved spacetimes. It also allows to formulate the theory of effective QG.



- AQFT is a convenient framework to solve conceptual problems of QFT on curved spacetimes. It also allows to formulate the theory of effective QG.
- In our framework, physical diffeomorphism invariant quantities can be viewed in 3 ways:



- AQFT is a convenient framework to solve conceptual problems of QFT on curved spacetimes. It also allows to formulate the theory of effective QG.
- In our framework, physical diffeomorphism invariant quantities can be viewed in 3 ways:
 - as locally covariant fields $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$,



- AQFT is a convenient framework to solve conceptual problems of QFT on curved spacetimes. It also allows to formulate the theory of effective QG.
- In our framework, physical diffeomorphism invariant quantities can be viewed in 3 ways:
 - as locally covariant fields $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$,
 - as covariant functionals $\Phi_f(g, X)$,



- AQFT is a convenient framework to solve conceptual problems of QFT on curved spacetimes. It also allows to formulate the theory of effective QG.
- In our framework, physical diffeomorphism invariant quantities can be viewed in 3 ways:
 - as locally covariant fields $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$,
 - as covariant functionals $\Phi_f(g, X)$,
 - as relational observables $\Phi_f(., X_g)$.



- AQFT is a convenient framework to solve conceptual problems of QFT on curved spacetimes. It also allows to formulate the theory of effective QG.
- In our framework, physical diffeomorphism invariant quantities can be viewed in 3 ways:
 - as locally covariant fields $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$,
 - as covariant functionals $\Phi_f(g, X)$,
 - as relational observables $\Phi_f(., X_g)$.
- To quantize the theory, we make a tentative split into a free and interacting part. We quantize the free theory first and then use the Epstein-Glaser renormalization to introduce the interaction.



- AQFT is a convenient framework to solve conceptual problems of QFT on curved spacetimes. It also allows to formulate the theory of effective QG.
- In our framework, physical diffeomorphism invariant quantities can be viewed in 3 ways:
 - as locally covariant fields $\Phi_{\mathcal{M}} : \mathfrak{D}(\mathcal{M}) \to \mathfrak{F}(\mathcal{M})$,
 - as covariant functionals $\Phi_f(g, X)$,
 - as relational observables $\Phi_f(., X_g)$.
- To quantize the theory, we make a tentative split into a free and interacting part. We quantize the free theory first and then use the Epstein-Glaser renormalization to introduce the interaction.
- We have shown that our theory is background independent, i.e. independent of the split into free and interacting part.





Thank you for your attention!