

Deformations of space-time metrics

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- To adapt a known exact solution to a more realistic one ...

Solutions with symmetries

Schwarzschild 's metric

$$ds^2 = - \left(1 - \frac{2MG}{c^2 r}\right) dt^2 + \left(1 - \frac{2MG}{c^2 r}\right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

FLRW Universe

$$ds^2 = c^2 dt^2 - a^2(t) \left[\frac{dr^2}{1 - kr^2} + r^2 d\Omega^2 \right]$$

$$g = \Omega \eta$$

Deformations 1

$$(M, g)$$

$$g_{ab} = \eta_{AB} \omega_a^A \omega_b^B, \quad g^{ab} = \eta^{AB} e_A^a e_B^b,$$

Tensorial
index

Tetrad
index

Matrix of scalar
field

$$\phi_B^A$$

$$\tilde{g}_{ab} \equiv \eta_{AB} \phi_C^A \phi_D^B \omega_a^C \omega_b^D$$

Deformed
tensor

• “First deforming matrix”

$$\psi_D^A \equiv \phi_D^{-1A}$$

$$\phi_B^A$$

$$\omega_a^A \phi_A^B \equiv \tilde{\omega}_a^B$$

$$\tilde{\omega}_a^A$$

$$\psi_B^A$$

$$\tilde{e}_A^a \equiv \psi_A^B e_B^a$$

$$\tilde{e}_A^a$$

$$\tilde{\omega}_a^A \tilde{e}_B^a = \delta_B^A, \quad \tilde{\omega}_a^A \tilde{e}_A^b = \delta_a^b.$$

Deformations 2

$$\tilde{g}_{ab} = \eta_{AB} \phi^A_C \phi^B_D \omega^C_a \omega^D_b$$

$$\tilde{g}^{ab} = \eta^{AB} \psi^C_A \psi^D_B e^a_C e^b_D$$

$$\tilde{g}_{ab} = \eta_{AB} \tilde{\omega}^C_a \tilde{\omega}^D_b$$

$$\tilde{g}^{ab} = \eta^{AB} \tilde{e}^a_C \tilde{e}^b_D$$

- “Second deforming matrix”

$$\mathcal{G}_{CD} \equiv \eta_{AB} \phi^A_C \phi^B_D, \quad \text{e} \quad \mathcal{G}^{CD} \equiv \eta_{AB} \psi^C_A \psi^D_B$$

$$\tilde{g}_{ab} = \mathcal{G}_{CD} \omega^C_a \omega^D_b, \quad \tilde{g}^{ab} = \mathcal{G}^{CD} e^a_C e^b_D,$$

Some features of deformations

•Commutation relations

$$[\tilde{e}_A, \tilde{e}_B]^b = \psi_A^S \psi_B^T [e_S, e_T]^b + 2\psi_{[A}^S e_{|S}^a e_{T]}^b (\partial_a \psi_{B]}^T)$$

•Identity

$$g_{ab} = \eta_{AB} \omega_a^A \omega_b^B, \quad g_{ab} = \eta_{AB} \omega_a'^A \omega_b'^B = \eta_{AB} \Lambda_C^A \Lambda_D^B \omega_a^C \omega_b^D$$

$$\Lambda^A_B \phi^B_C$$

•Diffeomorfism

$$\tilde{g}_{\alpha\beta} = g_{\mu\nu} \phi^\mu_\alpha \phi^\nu_\beta$$

$$\phi^a_b \equiv \phi^A_B \omega_b^B e_A^a$$

Deformations of tridimensional metrics

$$\textcircled{g} = \sigma h + \epsilon s \otimes s$$

Deformed
metric

Costant
curvature
metric

1--form

$$\begin{aligned} h_{ab} &= \eta_{AB} \omega_a^A \omega_b^B, \\ g_{ab} &= \mathcal{G}_{CD} \omega_a^C \omega_b^D \end{aligned}$$

$$\tilde{g}_{ab} = \eta_{AB} \phi^A_C \phi^B_D \omega_a^C \omega_b^D$$

$$\phi^A_C = \sqrt{\sigma} \delta^A_C + \alpha s^A s_C$$

$$f = \frac{n(n-1)}{2}.$$

$$\|s\|_h = \sigma,$$

Conformal deformations

$$\hat{g} \equiv \Omega^2 g = \Omega^2 \eta_{AB} \omega^A \otimes \omega^B$$

$$\psi^c{}_A = \Omega \Lambda^c{}_A$$

$$\tilde{g}_{ab} = g_{ab} + h_{ab}$$

Deformed
metric

$$\tilde{g}^{ab} = g^{ab} + X^{ab}$$

Tensor

$$X^{ab} = \mathcal{D}^{AB} e_A^a e_B^b$$

$$\mathcal{G}_{AB} = \eta_{AB} + \mathcal{C}_{AB}$$

$$\mathcal{G}^{AB} = \eta^{AB} + \mathcal{D}^{AB}$$

$$\mathcal{D}^{AD} = -\eta^{AB} \mathcal{G}^{CD} \mathcal{C}_{BC}$$

$$X^{ab} h_{bc} \neq \delta_c^a$$

$$X^{ad} = -g^{ab} \tilde{g}^{cd} h_{bc}$$

$$\mathcal{D}^{AB} \mathcal{C}_{BC} \neq \delta^A_C$$

$$\tilde{g}_{ab} = g_{ab} + h_{ab}, \quad \tilde{g}^{ab} = g^{ab} + X^{ab}$$

$$\mathcal{G}_{AB} = \eta_{CD} \phi^C_A \phi^D_B$$

$$\phi^A_B = \delta^A_B + \varphi^A_B$$

$$\varphi^A_A = 0, \quad \varphi^A_0 = \varphi^0_A, \quad \varphi^A_i = -\varphi^i_A \quad \text{dove} \quad i \in \{1, 2, 3\}$$

$$\varphi^B_A = \begin{pmatrix} 0 & a & b & c \\ a & 0 & d_1 & d_2 \\ b & -d_1 & 0 & d_3 \\ c & -d_2 & -d_3 & 0 \end{pmatrix}$$

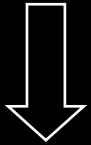
$$\begin{aligned} \tilde{g}_{ab} &= \Omega^2 g_{ab} + h_{ab} \\ \tilde{g}^{ab} &= \frac{1}{\Omega^2} g^{ab} + X^{ab} \end{aligned}$$

$$h_{ab} \equiv \eta_{AB} \varphi^A_C \varphi^B_D \omega^C_a \omega^D_b, \quad X^{ad} \equiv -\Omega^{-2} g^{ab} \tilde{g}^{cd} h_{bc}$$

$$\phi^A_B = \Omega \delta^A_B + \varphi^A_B; \quad \varphi = \begin{pmatrix} 0 & a & b & c \\ a & 0 & d_1 & d_2 \\ b & -d_1 & 0 & d_3 \\ c & -d_2 & -d_3 & 0 \end{pmatrix}$$

•Gravitational Waves

$$|\mathcal{C}_{AB}| \ll 1 \quad \forall A \in B,$$



$$\tilde{g}^{ab} = g^{ab} - h^{ab}$$

Gauge



$$\nabla^a h_{ab} = 0, \\ h = 0.$$

$$h \equiv g^{bd} h_{bd}$$

$$G_{ab} = 0$$

$$\tilde{g}_{ab} = g_{ab} + h_{ab}$$

$$h_{ab} = \mathcal{C}_{AB} \omega_a^A \omega_b^B$$

$$\mathcal{G}_{AB} = \eta_{AB} + \mathcal{C}_{AB}$$

Equations

$$\square h_{ac} = 2 \underline{R^b_{ac}}^d h_{bd}$$

$$(\square \mathcal{C}_{AB}) \omega_a^A \omega_b^B = \mathcal{C}_{AB} [2 R^b_{ac}{}^d \omega_b^A \omega_d^B - \square (\omega_a^A \omega_b^B)]$$

Deformations of
Minkowski space time

$$\square h_{ac} = 0$$



$$(\square \mathcal{C}_{AB}) \omega_a^A \omega_b^B = 0$$

- Geodesic motions

$$v^a \nabla_a v^b = 0$$

Gravitational lensing

Geodesic deviations

Gravitational redshift

$$v^a \tilde{\nabla}_a v^d = \frac{1}{2} \tilde{g}^{cd} \{ \nabla_b \tilde{g}_{ac} + \nabla_a \tilde{g}_{bc} - \nabla_c \tilde{g}_{ba} \} v^a v^b$$

$$v^a \tilde{\nabla}_a v^d = \frac{1}{2} \tilde{g}^{cd} \{ \nabla_a h_{cb} + \nabla_b h_{ca} - \nabla_c h_{ba} \} v^a v^b$$

$$\tilde{g}_{ab} = g_{ab} + h_{ab}, \quad \tilde{g}^{ab} = g^{ab} + X^{ab}$$

$$X^{dc} (h_{cb,a} + h_{ca,b} - h_{ba,c}) v^a v^b + X^{dc} (g_{cb,a} + g_{ca,b} - g_{ba,c}) v^a v^b + g^{dc} (h_{cb,a} + h_{ca,b} - h_{ba,c}) v^a v^b = 0.$$

$$\begin{aligned} \tilde{g}_{ab} &= \Omega^2 g_{ab} + h_{ab} \\ \tilde{g}^{ab} &= \frac{1}{\Omega^2} g^{ab} + X^{ab} \end{aligned}$$

$$\begin{aligned} v^a \tilde{C}_{ac}^b v^c &= 2(\ln \Omega)_{,v} v^b - (g_{ac} v^a v^c) [g^{bd} \nabla_d \ln \Omega + X^{bd} \nabla_d \ln \Omega] + \\ &+ 2v^a (\nabla_a \ln \Omega) g_{dc} X^{db} v^c + \frac{1}{2} v^a [(\nabla_a h_{cd}) + (\nabla_c h_{ad}) + \\ &- (\nabla_d h_{ac})] v^c (g^{bd} \Omega^{-2} + X^{bd}) \end{aligned}$$

- Gravitational lensing

$$\frac{dp^y}{d\lambda} = -\tilde{\Gamma}_{\beta\gamma}^y p^\beta p^\gamma$$

$$\Delta p^y \equiv - \int_{-\infty}^{+\infty} \Gamma_{\beta\gamma}^y p^\beta p^\gamma d\lambda,$$

$$\delta p^y = -\frac{1}{2} \int_{-\infty}^{+\infty} \tilde{g}^{y\delta} \left[\partial_\beta (\mathcal{G}_{\text{CD}} \omega_\gamma^{\text{C}} \omega_\delta^{\text{D}}) + \partial_\gamma (\mathcal{G}_{\text{CD}} \omega_\beta^{\text{C}} \omega_\delta^{\text{D}}) - \partial_\delta (\mathcal{G}_{\text{CD}} \omega_\beta^{\text{C}} \omega_\gamma^{\text{D}}) \right] p^\beta p^\gamma d\lambda,$$

$$\delta p^y = -\frac{1}{2} \eta_{\text{QP}} \int_{-\infty}^{+\infty} \tilde{g}^{y\delta} \left[\partial_\beta (\tilde{\omega}_\gamma^{\text{Q}} \tilde{\omega}_\delta^{\text{P}}) + \partial_\gamma (\tilde{\omega}_\beta^{\text{Q}} \tilde{\omega}_\delta^{\text{P}}) - \partial_\delta (\tilde{\omega}_\beta^{\text{Q}} \tilde{\omega}_\gamma^{\text{P}}) \right] p^\beta p^\gamma d\lambda .$$

$$\delta p^y = \Delta p^y - \frac{1}{2} \int_{-\infty}^{+\infty} [X^{y\delta} (g_{\delta\beta,\gamma} + g_{\delta\gamma,\beta} - g_{\beta\gamma,\delta}) + g^{y\delta} (h_{\delta\beta,\gamma} + h_{\delta\gamma,\beta} - h_{\beta\gamma,\delta})] p^\beta p^\gamma d\lambda +$$

$$- \frac{1}{2} \int_{-\infty}^{+\infty} X^{y\delta} (h_{\delta\beta,\gamma} + h_{\delta\gamma,\beta} - h_{\beta\gamma,\delta}) p^\beta p^\gamma d\lambda$$

- Gravitational lensing

$$\frac{d\mathbf{e}}{dl_{eucl}} = -\frac{2}{c^2} \nabla_{\perp} \Phi$$

$$ds^2 = \left(1 + \frac{2\Phi}{c^2}\right) c^2 dt^2 - \left(1 - \frac{2\Phi}{c^2}\right) \delta_{ij} dx^i dx^j$$

$$\nabla_{\perp} \Phi \equiv [\nabla \Phi - \mathbf{e}(\mathbf{e} \cdot \nabla \Phi)]$$

$$\tilde{ds}^2 = A(t) \left(1 + \frac{2\Phi}{c^2}\right) c^2 dt^2 - A(t) \left(1 - \frac{2\Phi}{c^2}\right) \delta_{ij} dx^i dx^j$$

$$\frac{d\mathbf{e}}{dl_{eucl}} = -\frac{2}{c^2} \nabla_{\perp} \Phi - (\partial_0 \ln A) \mathbf{e}$$

$$\tilde{ds}^2 = A(|x|) \left(1 + \frac{2\Phi}{c^2}\right) c^2 dt^2 - A(|x|) \left(1 - \frac{2\Phi}{c^2}\right) \delta_{ij} dx^i dx^j$$

$$\frac{d\mathbf{e}}{dl_{eucl}} = -\nabla_{\perp} \hat{\Phi}$$

$$\hat{\Phi} \equiv \frac{2}{c^2} \Phi + \ln A$$

$$|x|^2 = x^2 + y^2 + z^2$$

- Geodetic deviations

$$a^a = T^c \tilde{\nabla}_c v^a$$

$$a^a = -\tilde{R}_{cbd}{}^a S^b T^c T^d.$$

$$a_0^a \equiv -R_{cbd}{}^a S^b T^c T^d$$

$$a^a = a_0^a - \overline{R}_{cbd}{}^a S^b T^c T^d$$

$$a^a = -2(\partial_{[b} \tilde{C}^a_{|d|c]} + \tilde{C}^f_{[c|d]} \tilde{\Gamma}^a_{b]f} + \Gamma^f_{[c|d]} \tilde{C}^a_{b]f}) S^b T^c T^d$$

- Gravitational redshift

$$z \equiv \frac{\lambda_2 - \lambda_1}{\lambda_1}$$



$$z = \frac{\nu_1}{\nu_2} - 1 = \frac{\sqrt{-\underline{\mathcal{G}}_{AB} \omega_0^A \omega_0^B}|_{P_2}}{\sqrt{-\underline{\mathcal{G}}_{AB} \omega_0^A \omega_0^B}|_{P_1}} - 1$$

Deformations of space time metrics

Conclusion

We have seen it is possible to deform space--time metrics using a matrix of scalar fields defined on the manifold.

Deformations could be a generalization of the relation

$$g = \Omega \eta$$

Deformations of space-time metrics could be a method to “adapt” a model solution to a more realistic one.

Deformations could be a way to study the inverse problem in general relativity.

Deformations of space time metrics

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